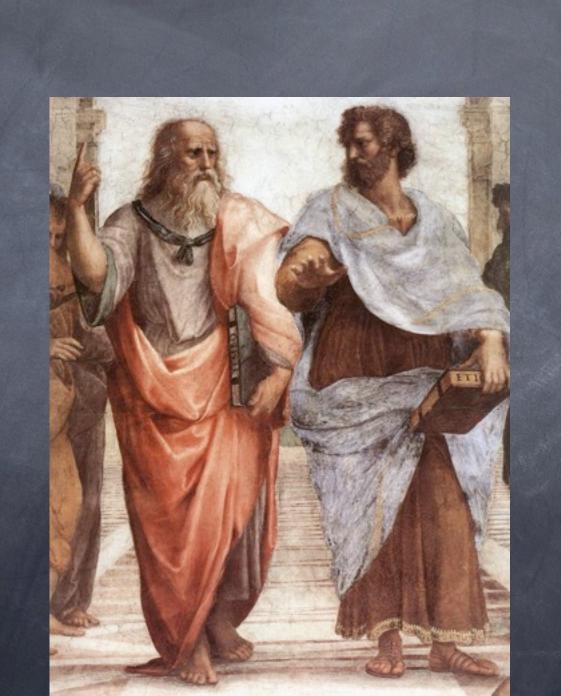
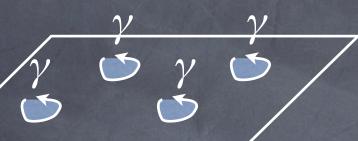
In search of topological states with half quantum vortices.

Eun-Ah Kim Cornell University

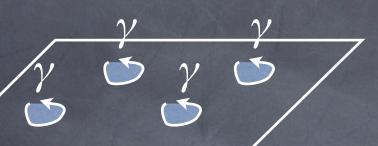
Suk-Bum Chung (Stanford)
Subroto Mukerjee (Berkeley)

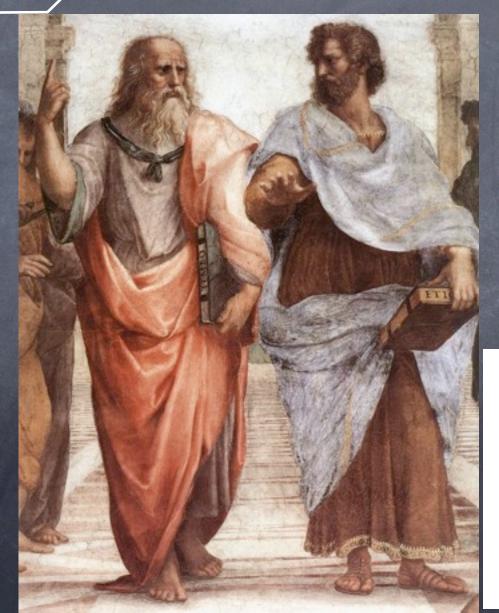
Hendrik Bluhm (Harvard)
Daniel Agterberg (UWM)

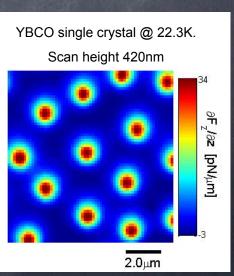






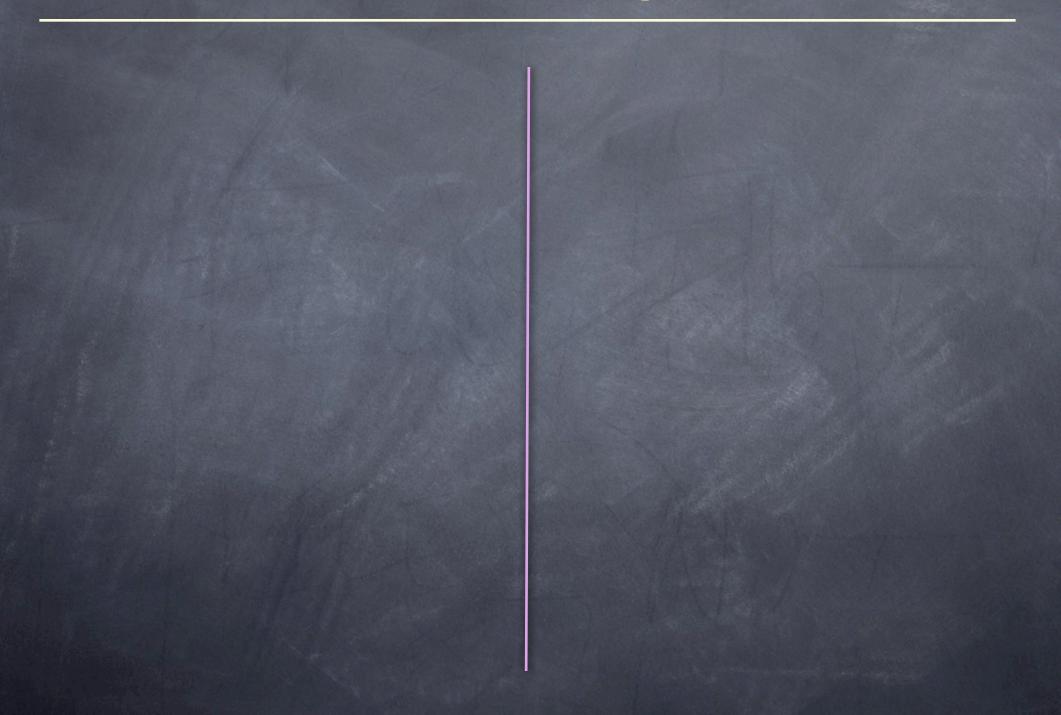






In search of topological states with half quantum vortices

- Topological order and fractionalization
- 1/2 QV's
- Stability of 1/2-QV's in SrRuO
- 1/2 QV lattices







Conventional order



Symmetry of the underlying Hamiltonian.



- Symmetry of the underlying Hamiltonian.
- reduced symmetry



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- Local measurements
 - → order parameter

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Gapped spectrum

No local order parameter.

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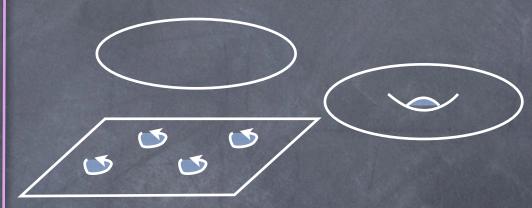
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Topological order



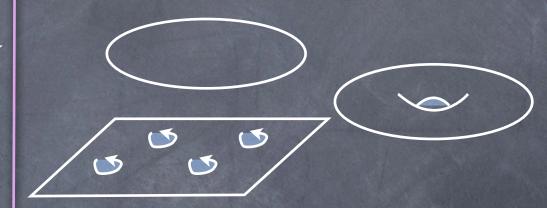
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Topological order



- Gapped spectrum
- Topological invariance
 - emergent symmetry
- No local order parameter.
- Topological degeneracy N_g.

Sweet Topology

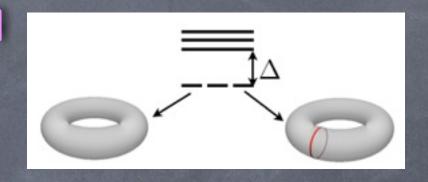
Sweet Topology



Sweet Topology

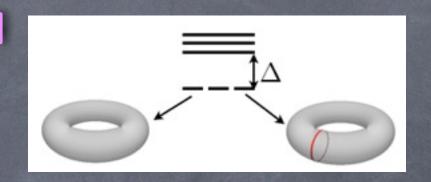


Fractional charge $e^*=e/q$ $N_g=q^g$ e.g., $N_1=3$



Wen and Niu, PRB, 1990 Stone and Chung, PRB, 2006

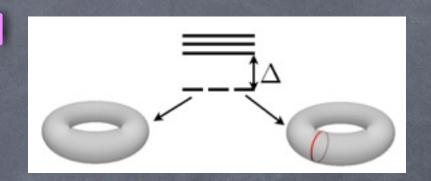
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2n Non-abelian vortices

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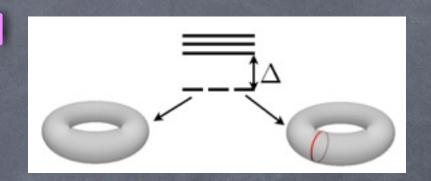


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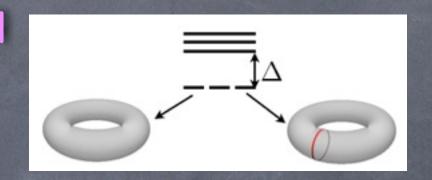
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 $N_{2n}=2^{n-1}$ for MR state or p+ip SF

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$$\Psi(x_1, \cdots, x_n) = c$$
-number

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exchange of qp's:

phase multiplication

to a complex number

$$\Psi(x_1 \leftrightarrow x_3) = e^{i\theta} \Psi$$

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A. Leggett, RMP 1975 Sigrist & Ueda RMP 1991

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© Gap function
$$\Delta_{ss'}(\mathbf{k}) = -\sum_{\mathbf{k}',s_3,s_4} V_{s'ss_3s_4}(\mathbf{k},\mathbf{k}') \langle a_{\mathbf{k}'s_3}a_{-\mathbf{k}'s_4} \rangle$$

$$\widehat{\Delta}(\mathbf{k}) = -\widehat{\Delta}^T(-\mathbf{k})$$

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Triplet gap matrix

$$\begin{split} \widehat{\Delta}(\mathbf{k}) &= i(\mathbf{d}(\mathbf{k}) \cdot \widehat{\boldsymbol{\sigma}}) \widehat{\boldsymbol{\sigma}}_{y} \\ &= \begin{bmatrix} -d_{x}(\mathbf{k}) + id_{y}(\mathbf{k}) & d_{z}(\mathbf{k}) \\ d_{z}(\mathbf{k}) & d_{x}(\mathbf{k}) + id_{y}(\mathbf{k}) \end{bmatrix} \end{split}$$

p+ip SC

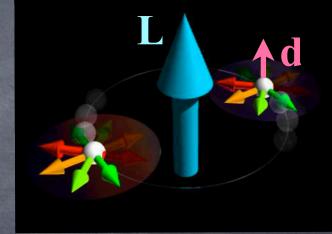
T-breaking (ABM)

$$\Delta(\hat{\mathbf{k}}) = \Delta_0(T)(p_x \pm ip_y) \begin{pmatrix} -d_x + id_y & d_z \\ d_z & d_x + id_y \end{pmatrix}$$

where $\hat{\mathbf{d}}$ is a real unit vector

In plane \mathbf{d} $d_z=0 \text{ i.e., } \mathbf{d}=(\cos\alpha,\sin\alpha,0)$

$$\Delta(\hat{\mathbf{k}}) = \Delta_0(T)(p_x \pm ip_y) \begin{pmatrix} -e^{-i\alpha} & 0\\ 0 & e^{i\alpha} \end{pmatrix}$$



1/2 QV with in-plane â

1/2 QV with in-plane à

The gap matrix

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- ② 1/2 QV when $d = (\cos \alpha, \sin \alpha, \theta)$
 - \leftrightarrow 2π winding for only one spin component
 - \rightarrow π winding of order parameter phase ϕ + π rotation of d vector

$$\Delta \alpha = \pm \pi$$
 $\Delta \phi = 2\pi$
 $\Delta \phi = \pi$
 $hc/2e$ vortex $hc/4e$ vortices

ØVortices of p+ip SF → zero modes at the core

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$$\gamma$$
 γ γ

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Das Sarma, Tewari, Nayak (06) Stone & Chung(06) Ivanov(01)

1/2 QV's: single Majorana zero mode

5/2 state described as p+ip paired stated of composite fermion

Pfaffian is real space many body BCS wave function of p+ip SF

HQV is equivalent to 1/4 qp

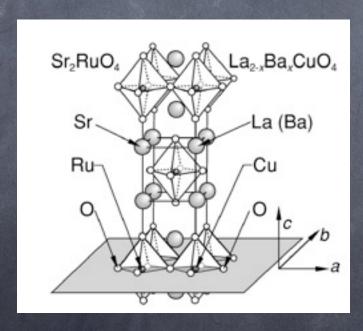
Moore & Read (91) Read & Green (00) Schriffer, p 48

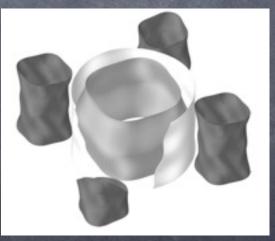
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K. Ishida et al, Nature (1998)

Spin-triplet superconductivity in Sr₂RuO₄ identified by ¹⁷O Knight shift





Experiments?

- NMR on ³He-A thin films: X Hakonen et al. Physica (89)
- Small angle neutron scattering: X Riseman et al. Nature (98)
- Scanning SQUID imaging: X

 Dolocan et al, PRL (05), Bjorsson et al, PRB (05)
- **O** NMR in the presence of $\mathbf{H} \perp ab$
 - ▶ d // ab: for $H_{\perp} \approx 200 \; G$, Murakawa et al, PRL (04)

Energy competition between full-QV and 1/2-QV

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$$f_{\rm grad}^{\rm 2D} = \frac{1}{2} \left(\frac{\hbar}{2m}\right)^2 \!\! \left[\rho_{\rm s} \! \left(\! \nabla_{\!\! \perp} \! \phi \! - \! \frac{2e}{\hbar c} \mathbf{A} \! \right)^2 \!\! + \! \rho_{\rm sp} \left(\! \nabla_{\!\! \perp} \! \alpha \right)^2 \right] + \frac{1}{8\pi} (\nabla \times \mathbf{A})^2$$

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Spin current energy diverges logarithmically!

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Spin current energy diverges logarithmically!

$$\epsilon_{\rm sp} = \frac{\pi}{4} \left(\frac{\hbar}{2m}\right)^2 \rho_{\rm sp} \ln \left(\frac{R}{\xi}\right)$$

stability of 1/2 QV



$$E_{\text{pair}}^{\text{half}}(r_{12}) = \frac{1}{2} \frac{\Phi_0^2}{16\pi^2 \lambda^2} \left[\ln\left(\frac{\lambda}{\xi}\right) + K_0\left(\frac{r_{12}}{\lambda}\right) + \frac{\rho_{\text{sp}}}{\rho_{\text{s}}} \ln\left(\frac{r_{12}}{\xi}\right) \right]$$

$$\Delta \chi = 2\pi$$

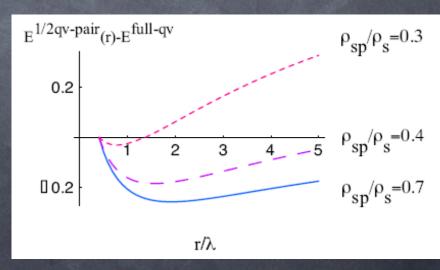
$$E^{\text{full}} = \pi \left(\frac{\hbar}{2m}\right)^2 \rho_{\text{s}} \ln \left(\frac{\lambda}{\xi}\right) = \frac{\Phi_0^2}{16\pi^2 \lambda^2} \ln \left(\frac{\lambda}{\xi}\right)$$

Competition between screened magnetic repulsion

and unscreened spin attraction

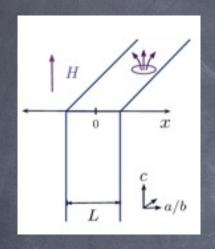
Finite equilibrium size for small ρ_{sp}/ρ_s

Leggett RMP 75

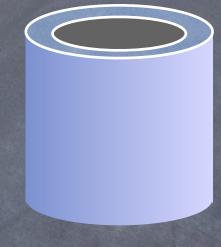


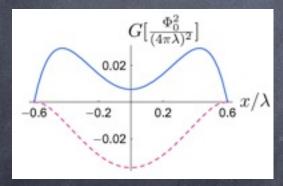
Mesoscopic sample

 \circ Sample of size $\sim \lambda$ a few micron



$$L=2\lambda$$





Underway in Budakian lab

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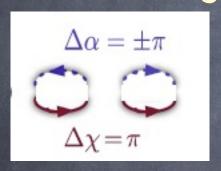
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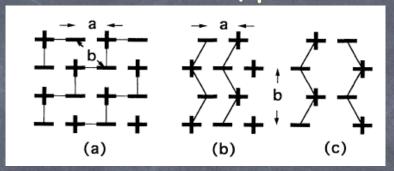
Natural way to stabilize 1/2 QV

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 - -T. Riseman at el , Nature (98) confirmed square lattice

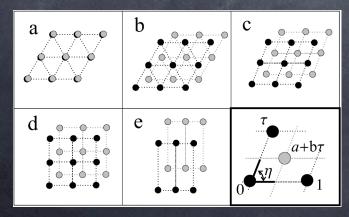
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- Potential of tuning $ho_{sp}/
 ho_s$
 - -Knowledge exist for $ho_{sp}/
 ho_s$ as a function of Fermi liquid parameters
 - -p-wave Feshbach resonance

- \odot SC (SF) with additional U(1) symmetry due to $\hat{\mathbf{d}}$ rotation
- Interlacing lattices of two types of vortices





- -Different geometry depending on density and LL mixing
- Similar case arise in spinor condensate

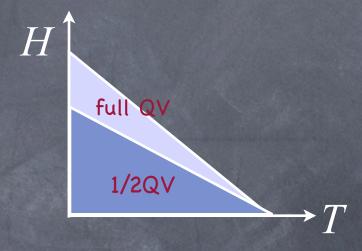


Muller & Ho(02)

Barnett, Mukergee & Moore(08)

Prediction

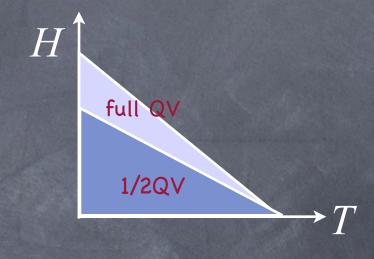
- Minimze GL free energy to determine the VL structure
- Quartic terms in the free energy determine the structure

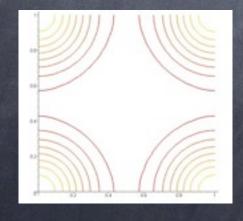


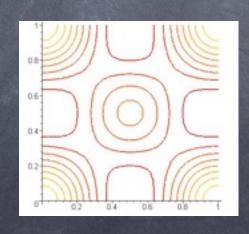
Field distribution as can be measured by neutron

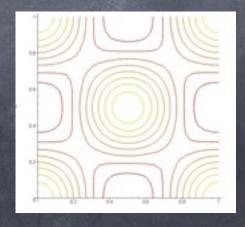
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Field distribution as can be measured by neutron

Stiffness engineering?

p-wave Feshbach resonance can allow for tuning for Fermi liquid parameters

$$H = \sum_{\mathbf{p}} \epsilon(\mathbf{p}) a_{\mathbf{p}}^{\dagger} a_{\mathbf{p}} + \sum_{\mathbf{p},\alpha} \left[\epsilon_{\alpha} + \frac{\epsilon(\mathbf{p})}{2} \right] b_{\mathbf{p}\alpha}^{\dagger} b_{\mathbf{p}\alpha} + \frac{1}{\sqrt{V}} \sum_{\mathbf{p},\mathbf{q},\alpha} g_{\mathbf{p}} p_{\alpha} \left(b_{\mathbf{q}\alpha} a_{\mathbf{p} + \frac{\mathbf{q}}{2}}^{\dagger} a_{-\mathbf{p} + \frac{\mathbf{q}}{2}}^{\dagger} + \text{h.c.} \right)$$

Gurarie, L. Radzihovsky, & A. V. Andreev (05)





@1/2 QV's are not stable in bulk systems



- 1/2 QV's are not stable in bulk systems
- Mesoscopic samples could favor 1/2 QV's



- 1/2 QV's are not stable in bulk systems
- Mesoscopic samples could favor 1/2 QV's
- @1/2 QV Vortex Lattice can be pursued and detected