



Eun-Ah Kim Cornell University

Acknowledgements



Prof. Michael Lawler Binghamton, Cornell



Dr. Kai Sun UIUC →UMD







Prof. James SethnaDr. Andy SchmitProf. Seamus DavisCornellCornellCornell, BNL

• Introduction

YBCO: spin and charge
BSCCO: got nematic?

• Look out

Reinitzer(1888)







Crystal

Т

Reinitzer(1888)



Crystal

Т

Liquid

Nematic



Liquid

Nematic

Smectic

Crystal



Reinitzer(1888)



Reinitzer(1888)





Sufficiently correlated systems





Pnictides

License to exist?

• Quantifying broken symmetry

License to exist?

• Quantifying broken symmetry

> Phenomenological model

License to exist?

• Quantifying broken symmetry

Phenomenological model

> Why would cuprates do that?

License to exist?

• Quantifying broken symmetry

> Phenomenological model

> Why would cuprates do that?

License to exist?

• Quantifying broken symmetry

Phenomenological model

> Why would cuprates do that?

License to exist?

Quantifying broken symmetry

Phenomenological model

Why would cuprates do that?

Nematic QCP, license to exist?

PHYSICAL REVIEW B 77, 184514 (2008)

Theory of the nodal nematic quantum phase transition in superconductors

Eun-Ah Kim,¹ Michael J. Lawler,² Paul Oreto,¹ Subir Sachdev,³ Eduardo Fradkin,³ and Steven A. Kivelson¹ ¹Department of Physics, Stanford University, Stanford, California 94305, USA ²Department of Physics, University of Toronto, Toronto, Ontario, Canada ³Department of Physics, University of Illinois at Urbana-Champaign, 1110 West Green Street, Urbana, Illinois 61801-3080, USA (Received 15 February 2008; published 22 May 2008)

 Nodal nematic QCP deep inside d-wave SC

Nematic QCP, license to exist?

PHYSICAL REVIEW B 77, 184514 (2008)

Theory of the nodal nematic quantum phase transition in superconductors

 Eun-Ah Kim,¹ Michael J. Lawler,² Paul Oreto,¹ Subir Sachdev,³ Eduardo Fradkin,³ and Steven A. Kivelson¹
 ¹Department of Physics, Stanford University, Stanford, California 94305, USA
 ²Department of Physics, University of Toronto, Toronto, Ontario, Canada
 ³Department of Physics, University of Illinois at Urbana-Champaign, 1110 West Green Street, Urbana, Illinois 61801-3080, USA (Received 15 February 2008; published 22 May 2008)

• Nodal nematic QCP deep inside d-wave SC



nematic QCP inside SC phase?

Nematic QCP, license to exist?

PHYSICAL REVIEW B 77, 184514 (2008)

Theory of the nodal nematic quantum phase transition in superconductors

Eun-Ah Kim,¹ Michael J. Lawler,² Paul Oreto,¹ Subir Sachdev,³ Eduardo Fradkin,³ and Steven A. Kivelson¹ ¹Department of Physics, Stanford University, Stanford, California 94305, USA ²Department of Physics, University of Toronto, Toronto, Ontario, Canada ³Department of Physics, University of Illinois at Urbana-Champaign, 1110 West Green Street, Urbana, Illinois 61801-3080, USA (Received 15 February 2008; published 22 May 2008)

• Nodal nematic QCP deep inside d-wave SC

• Nematic d-SC: d-SC + small s-component $\Delta_d (\cos k_x - \cos k_y) + \lambda \phi$



nematic QCP inside SC phase?

Looking for nematic critical fluctuations

•Self energy $\hat{\Sigma}(\vec{q}, \omega)$ due to fluctuation :k-selective decoherence



Interference of nematic quantum critical quasiparticles: a route to the octet model

Eun-Ah Kim¹ and Michael J. Lawler^{2, 1}

¹Department of Physics, Cornell University, Ithaca, NY 14853 ²Department of Physics, Binghamton University, Binghamton NY 13902 (Dated: November 13, 2008)

arXiv:0811.2242



YBCO, spin and charge





Anisotropy in inelastic $(\pi/a,\pi/b)$ neutron scattering peak Hinkov et al, Science, Jan 2008





Anisotropy in inelastic $(\pi/a,\pi/b)$ neutron scattering peak Hinkov et al, Science, Jan 2008





Anisotropy in inelastic $(\pi/a,\pi/b)$ neutron scattering peak Hinkov et al, Science, Jan 2008





ð

Anisotropy in inelastic $(\pi/a,\pi/b)$ neutron scattering peak Hinkov et al, Science, Jan 2008

Incommensurability as Order parameter?

 $\delta \propto \sqrt{T - T^*}$





Spin channel

Charge channel

Kai Sun,^{1,2} Michael J. Lawler,^{3,4} and Eun-Ah Kim⁴



Kai Sun,^{1,2} Michael J. Lawler,^{3,4} and Eun-Ah Kim^4





Kai Sun,^{1,2} Michael J. Lawler,^{3,4} and Eun-Ah Kim^4



$$S[\vec{\phi}] = \frac{1}{g} \int \frac{d^2 \mathbf{q} d\omega}{(2\pi)^3} \left(i\Gamma |\omega| + \omega^2 - \Delta^2 \left(\mathbf{q}\right) \right) |\vec{\phi}(\mathbf{q},\omega)|^2$$



Kai Sun,^{1,2} Michael J. Lawler,^{3,4} and Eun-Ah Kim^4



$$S[\vec{\phi}] = \frac{1}{g} \int \frac{d^2 \mathbf{q} d\omega}{(2\pi)^3} \left(i\Gamma |\omega| + \omega^2 - \Delta^2 \left(\mathbf{q}\right) \right) |\vec{\phi}(\mathbf{q},\omega)|^2$$

$$\chi''(\omega, \mathbf{q}) = g \frac{\Gamma \omega}{(\Gamma \omega)^2 + (\omega^2 - \Delta(\mathbf{q})^2)^2}$$



Kai Sun,^{1,2} Michael J. Lawler,^{3,4} and Eun-Ah Kim⁴



$$S[\vec{\phi}] = \frac{1}{g} \int \frac{d^2 \mathbf{q} d\omega}{(2\pi)^3} \left(i\Gamma |\omega| + \omega^2 - \Delta^2 \left(\mathbf{q}\right) \right) |\vec{\phi}(\mathbf{q},\omega)|^2$$

$$\chi''(\omega, \mathbf{q}) = g \frac{\Gamma \omega}{(\Gamma \omega)^2 + (\omega^2 - \Delta(\mathbf{q})^2)^2}$$

$$\frac{1}{4\pi/5} \pi \frac{6\pi/5}{k}$$

$$\Delta^2(\mathbf{q};N) = \Delta_0^2(N) + c_0^2(N)q^2 - c_2^2(N)N(q_x^2 - q_y^2) + \cdots$$

Kai Sun,^{1,2} Michael J. Lawler,^{3,4} and Eun-Ah Kim^4



$$S[\vec{\phi}] = \frac{1}{g} \int \frac{d^2 \mathbf{q} d\omega}{(2\pi)^3} \left(i\Gamma |\omega| + \omega^2 - \Delta^2 \left(\mathbf{q}\right) \right) |\vec{\phi}(\mathbf{q},\omega)|^2$$

$$\chi''(\omega, \mathbf{q}) = g \frac{\Gamma \omega}{(\Gamma \omega)^2 + (\omega^2 - \Delta(\mathbf{q})^2)^2}$$



$$\Delta^2(\mathbf{q};N) = \Delta_0^2(N) + c_0^2(N)q^2 - c_2^2(N)N(q_x^2 - q_y^2) + \cdots$$

Kai Sun,^{1,2} Michael J. Lawler,^{3,4} and Eun-Ah Kim^4

arXiv:0906.3460



 $\Delta^2(\mathbf{q};N) = \Delta_0^2(N) + c_0^2(N)q^2 - c_2^2(N)N(q_x^2 - q_y^2) + c_0^2(N)q^2 - c_2^2(N)N(q_x^2 - q_y^2) + c_0^2(N)Q^2 - c_0^2(N)N(q_x^2 - q_y^2) + c_0^2(N)Q^2 - c_0^2(N)N(q_x^2 - q_y^2) + c_0^2(N)N(q_x^2 - q$

$$S[\vec{\phi}] = \frac{1}{g} \int \frac{d^2 \mathbf{q} d\omega}{(2\pi)^3} \left(i\Gamma |\omega| + \omega^2 - \Delta^2 \left(\mathbf{q}\right) \right) |\vec{\phi}(\mathbf{q},\omega)|^2$$

$$\chi''(\omega, \mathbf{q}) = g \frac{\Gamma \omega}{(\Gamma \omega)^2 + (\omega^2 - \Delta(\mathbf{q})^2)^2}$$



Kai Sun,^{1,2} Michael J. Lawler,^{3,4} and Eun-Ah Kim⁴



$$S[\vec{\phi}] = \frac{1}{g} \int \frac{d^2 \mathbf{q} d\omega}{(2\pi)^3} \left(i\Gamma |\omega| + \omega^2 - \Delta^2 \left(\mathbf{q}\right) \right) |\vec{\phi}(\mathbf{q},\omega)|^2$$

$$\chi''(\omega, \mathbf{q}) = g \frac{\Gamma \omega}{(\Gamma \omega)^2 + (\omega^2 - \Delta(\mathbf{q})^2)^2}$$

$$\Delta^2(\mathbf{q};N) = \Delta_0^2(N) + c_0^2(N)q^2 - c_2^2(N)N(q_x^2 - q_y^2) + \cdots$$

$$\delta \propto \sqrt{T - T^*}$$



Kai Sun,^{1,2} Michael J. Lawler,^{3,4} and Eun-Ah Kim^4



$$S[\vec{\phi}] = \frac{1}{g} \int \frac{d^2 \mathbf{q} d\omega}{(2\pi)^3} \left(i\Gamma |\omega| + \omega^2 - \Delta^2 \left(\mathbf{q}\right) \right) |\vec{\phi}(\mathbf{q},\omega)|^2$$

$$\chi''(\omega, \mathbf{q}) = g \frac{\Gamma \omega}{(\Gamma \omega)^2 + (\omega^2 - \Delta(\mathbf{q})^2)^2}$$

$$\Delta^2(\mathbf{q};N) = \Delta_0^2(N) + c_0^2(N)q^2 - c_2^2(N)N(q_x^2 - q_y^2) + \cdots$$

$$\delta \propto \sqrt{T - T^*}$$



BSCCO, got nematic?



$\begin{array}{c} dI/dV(\omega)\text{-map}\\ \text{McElroy et al, Nature 422, 592 (2003)}\\ \text{OD }T_c\text{=}86\text{K}\ (\text{p=}\) \end{array}$



 $\begin{array}{c} \mbox{R-map} \\ \mbox{Kohsaka et al, Science 315, 1380 (2007)} \\ \mbox{UD }T_c = 45 \mbox{K (p=0.08)} \end{array}$

a. e=0.4 **b.** e=0.6 1.8 0.69

Figure S7 a-f. A series of images displaying the real space conductance ratio *Z* as a function of energy rescaled to the local psuedogap value, $e = E/\Delta_1(\mathbf{r})$. Each pixel location was rescaled independently of the others. The common color scale illustrates that the bond centered pattern appears strongest in *Z* exactly at $E = \Delta_1(\mathbf{r})$.

 $\begin{array}{c} \text{Z-map}(\omega)\\ \text{Kohsaka et al, Nature 454, 1072 (2008)}\\ \text{UD }T_c=\!45\text{K} \end{array}$

HAMLET: Do you see yonder cloud that's almost in shape of a camel?

POLONIUS: By th'mass, and 'tis like a camel indeed.

HAMLET: Methinks it is like a weasel.

POLONIUS: It is backed like a weasel.

--W. Shakespeare (S. Chakravarty's perspectives Science 08)



HAMLET: Do you see yonder cloud that's almost in shape of a camel?

POLONIUS: By th'mass, and 'tis like a camel indeed.

HAMLET: Methinks it is like a weasel.

POLONIUS: It is backed like a weasel.

--W. Shakespeare (S. Chakravarty's perspectives Science 08)



M. Lawler et al, in prep. Challenge: An objective measure

Candidate broken symmetries

2.7

2.1

1.8

1.5

1.2

0.9

0.6

0.3

Z-map intensity at E = 150.0 meV



• Translational symmetry \hat{T}_a, \hat{T}_b • Rotational symmetry $\hat{R}_{\pi/2}$

Can we separately measure?

Need a \hat{T}_a, \hat{T}_b preserving order parameter

On the shoulder of

• Relating asymmetry to a quantitative measure $Z(\mathbf{r}, \mathbf{w}) = R(\mathbf{r})$

P. Anderson, N.P. Ong J. Phys. Chem. Solids, **67**,1(1993)

M.B.J. Meinders, H. Eskes, G.A. Sawatzky Phys. Rev. B, **48**, 3916 (1993)

M. Randeria et al, PRL **95**, 137001 (2005)

• Fourier filtering to look for stripe

$$N_f(\mathbf{r}, E) = \int d\mathbf{r}' f(\mathbf{r} - \mathbf{r}') N(\mathbf{r}', E),$$

$$f(\mathbf{r}) \propto \Lambda^2 e^{-r^2 \Lambda^2/2} [\cos(\pi x/2a) + \cos(\pi y/2a)].$$

C. Howald et al, S. Kivelson et al, PRB **67**, 014533 (2003) RMP **75**, 1201 (2003)











- Cex

 S_x

Z-map intensity at E = 150.0 meV



Listen to Bragg Peaks



• Bragg peak $\tilde{Z}(\vec{Q}_x) = \frac{1}{\sqrt{N}} \sum_{\vec{R}+\vec{d}} Z(\vec{R}+\vec{d})e^{-i\vec{Q}_x\cdot\vec{d}}$ $\vec{Q}_x = (2\pi/a, 0)$

• Need O sites $\tilde{Z}(\vec{Q}_x) = \bar{Z}_{Cu} - \bar{Z}_{O_x} + \bar{Z}_{O_y}, \quad \tilde{Z}(\vec{Q}_y) = \bar{Z}_{Cu} + \bar{Z}_{O_x} - \bar{Z}_{O_y}$ $\mathcal{O}_N \propto (\bar{Z}_{O_x} - \bar{Z}_{O_y})$





• Need O sites $\tilde{Z}(\vec{Q}_x) = \bar{Z}_{Cu} - \bar{Z}_{O_x} + \bar{Z}_{O_y}, \quad \tilde{Z}(\vec{Q}_y) = \bar{Z}_{Cu} + \bar{Z}_{O_x} - \bar{Z}_{O_y}$ $\mathcal{O}_N \propto (\bar{Z}_{O_x} - \bar{Z}_{O_y})$





• Need O sites $\tilde{Z}(\vec{Q}_x) = \bar{Z}_{Cu} - \bar{Z}_{O_x} + \bar{Z}_{O_y}, \quad \tilde{Z}(\vec{Q}_y) = \bar{Z}_{Cu} + \bar{Z}_{O_x} - \bar{Z}_{O_y}$ $\mathcal{O}_N \propto (\bar{Z}_{O_x} - \bar{Z}_{O_y})$









Nematic ordering in UD 45

 Q_{X}

 \bigcirc

Q_y



Extracted from published data, T=4K

Kohsaka et al, Nature 454, 1072 (2008)

Domain size in Z-map

Domain size in Z-map





Nematic domains

• Shift Q_x, Q_y to origin ("tune to the channel")

• Low pass filter (long distance physics)



Nematic domains

• Shift Q_x, Q_y to origin ("tune to the channel")

• Low pass filter (long distance physics)





Listen to channel S

Oriented stripe domains

• Shift S_x, S_y to origin ("tune to the channel")

 \mathbf{Q}_{V}

 $Q_{\rm X}$

• Low pass filter (long distance physics)



VOLUME 66, NUMBER 24

PHYSICAL REVIEW LETTERS

17 JUNE 1991

Weak Pinning and Hexatic Order in a Doped Two-Dimensional Charge-Density-Wave System

Hongjie Dai, Huifen Chen, and Charles M. Lieber Departments of Chemistry and Applied Physics, Columbia University, New York, New York 10027 (Received 11 July 1990; revised manuscript received 25 February 1991)

Example: VOLUME 66, NUMBER 24

PHYSICAL REVIEW LETTERS

17 JUNE 1991

Weak Pinning and Hexatic Order in a Doped Two-Dimensional Charge-Density-Wave System

Hongjie Dai, Huifen Chen, and Charles M. Lieber Departments of Chemistry and Applied Physics, Columbia University, New York, New York 10027 (Received 11 July 1990; revised manuscript received 25 February 1991)

Example: VOLUME 66, NUMBER 24

PHYSICAL REVIEW LETTERS

17 JUNE 1991

Weak Pinning and Hexatic Order in a Doped Two-Dimensional Charge-Density-Wave System

Hongjie Dai, Huifen Chen, and Charles M. Lieber

Departments of Chemistry and Applied Physics, Columbia University, New York, New York 10027 (Received 11 July 1990; revised manuscript received 25 February 1991)

Scanning-tunneling microscopy has been used to characterize the effects of Nb impurities on the incommensurate charge-density-wave (CDW) phase in 1T-TaS₂. Real- and reciprocal-space data indicate that disorder in the CDW is due to dislocations and small random rotations of the CDW The dislocations destroy translational order; however, calculations show that the orientational order is long range.

Example: VOLUME 66, NUMBER 24

PHYSICAL REVIEW LETTERS

17 JUNE 1991

Weak Pinning and Hexatic Order in a Doped Two-Dimensional Charge-Density-Wave System

Hongjie Dai, Huifen Chen, and Charles M. Lieber

Departments of Chemistry and Applied Physics, Columbia University, New York, New York 10027 (Received 11 July 1990; revised manuscript received 25 February 1991)

Scanning-tunneling microscopy has been used to characterize the effects of Nb impurities on the incommensurate charge-density-wave (CDW) phase in 1T-TaS₂. Real- and reciprocal-space data indicate that disorder in the CDW is due to dislocations and small random rotations of the CDW The dislocations destroy translational order; however, calculations show that the orientational order is long range.



Example: VOLUME 66, NUMBER 24

PHYSICAL REVIEW LETTERS

17 JUNE 1991

Weak Pinning and Hexatic Order in a Doped Two-Dimensional Charge-Density-Wave System

Hongjie Dai, Huifen Chen, and Charles M. Lieber

Departments of Chemistry and Applied Physics, Columbia University, New York, New York 10027 (Received 11 July 1990; revised manuscript received 25 February 1991)

Scanning-tunneling microscopy has been used to characterize the effects of Nb impurities on the incommensurate charge-density-wave (CDW) phase in 1T-TaS₂. Real- and reciprocal-space data indicate that disorder in the CDW is due to dislocations and small random rotations of the CDW The dislocations destroy translational order; however, calculations show that the orientational order is long range.



in the Nb-doped samples. These small rotations are readily observed by viewing the images at a glancing angle along the indicated lines. Analyses of atomic-

Example: VOLUME 66, NUMBER 24

PHYSICAL REVIEW LETTERS

17 JUNE 1991

Weak Pinning and Hexatic Order in a Doped Two-Dimensional Charge-Density-Wave System

Hongjie Dai, Huifen Chen, and Charles M. Lieber

Departments of Chemistry and Applied Physics, Columbia University, New York, New York 10027 (Received 11 July 1990; revised manuscript received 25 February 1991)

Scanning-tunneling microscopy has been used to characterize the effects of Nb impurities on the incommensurate charge-density-wave (CDW) phase in 1T-TaS₂. Real- and reciprocal-space data indicate that disorder in the CDW is due to dislocations and small random rotations of the CDW The dislocations destroy translational order; however, calculations show that the orientational order is long range.



in the Nb-doped samples. These small rotations are readily observed by viewing the images at a glancing angle along the indicated lines. Analyses of atomicElectronic Nematic in Cuprates? Looking ahead

- Doping dependence?
- Temperature dependence?
- Diffraction measurements?



Phenomenological model

> Why would cuprates do that?